In-situ sampling uncovers the dynamics of particle genesis and growth in an aerosol tube reactor

Eric Bain Wasmund · Kenneth S. Coley

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Abstract Making metal and ceramic powders using aerosol synthesis from vapour precursors, either in a flame or hot-wall tube reactor, is the basis for producing many thousands of tons of powder on an annual basis. To properly study this system, we have designed and built a model reactor with sampling points at evenly spaced axial positions. This allows us to take snapshots of the aerosol population at many points within the reactor. Nucleation followed by a surface reaction produces a solid phase extremely rapidly, within 0.01 s under typical conditions. This is followed by a transient state where nucleation, surface reaction and coagulation all interact to produce a strongly bimodal size distribution. After nucleation is extinguished, the size distribution approaches the selfpreserving limit as predicted for a coagulation-dominated process. The final structure is determined by the dominant sintering mechanism, which can be estimated from theory. The knowledge of this mechanism offers the possibility of selecting reactor conditions to produce powders with optimized properties.

More than 100 years ago, Ludwig Mond and Carl Langer discovered that a number of transition metals react reversibly with carbon monoxide to form a class of compounds known as the metal carbonyls [\[1](#page-6-0)]. In these

E. B. Wasmund (\boxtimes)

K. S. Coley

compounds, CO molecules are bound as ligands to the metal. They typically exist as vapour under ambient conditions, and are generally quite toxic. The equilibrium for dissociation of nickel carbonyl is shown as Eq. (1) [\[2](#page-6-0)].

$$
Ni (CO)4 \leftrightarrow Ni (s) + 4 CO
$$

\n
$$
\Delta G^{0}(kJ/mol) = 155.20 - 0.3964 T(K)
$$
\n(1)

From Eq. (1), nickel carbonyl is favoured by lower temperature and higher pressure. The discovery of these compounds paved the way for industrial innovations in the areas of extractive metallurgy and chemical vapour depo-sition [\[3](#page-6-0)].

Regarding extractive metallurgy, industrially significant processes for the low-cost purification of nickel by the carbonyl process have been in continuous operation since 1902. In these processes, metals that have the tendency to form carbonyls are readily separated from other metals by passing carbon monoxide through a bed of the finely divided feed, under conditions that favour the left-hand side of Eq. (1). Metals such as nickel will selectively form gaseous nickel carbonyl, leaving other metals such as copper, behind in the bed. The nickel carbonyl can then be deposited as pure nickel by raising the temperature.

Regarding chemical vapour deposition, metal carbonyls can be been used as chemical vapour deposition (CVD) precursors to make a variety of materials such as nickelcoated fibres, foam and particulates [\[4](#page-6-0)]. In these processes, metals are deposited heterogeneously from the gas-phase onto a heated template, again using the reaction described in Eq. (1). The kinetics for the heterogeneous reaction have been well studied and can be accurately described by a second order mechanism [[5\]](#page-6-0). As a result, an adequate

Energy Materials, Inco Special Products, 2101 Hadwen Rd, L5K 2L3 Mississauga, ON, Canada e-mail: ewasmund@inco.com

Department of Materials Science and Engineering, McMaster University, 1280 Main St W, L8S 4L7 Hamilton, ON, Canada

application of chemical kinetics, heat transfer and mass transport allows practitioners to reliably deposit metal coatings with precise control of thickness and microstructure.

The metal carbonyls can also be used to make powders by free-space decomposition. Iron and nickel powders have been made by this process since the 1920s using flowthrough tube reactors developed originally by Alwyn Mittasch [[3\]](#page-6-0). In these reactors, the metal carbonyl, in a suitable carrier gas such as carbon monoxide is fed into the top of a tube, typically 1–2 m in diameter and 5–10 m in height. The tube is heated through the walls, and as the metal carbonyl descends through the reactor it reaches a critical temperature, which initiates the decomposition reaction, resulting in the nucleation of a solid phase and liberation of CO. These particles continue to grow by a surface reaction as they are carried by gravity and gas entrainment to the bottom of the reactor where they are collected in a hopper.

The science that defines the control of particle properties in this process is much more complicated than chemical vapour deposition alone. Equation (1) can proceed by two parallel pathways, a homogeneous decomposition, or nucleation reaction, and a heterogeneous reaction on the surface of pre-existing particles. As these reactions occur, the particles can also grow by collision. This is known as coagulation in the aerosol literature, and the final particle structure can also be affected by gas-phase sintering. Besides metal carbonyl decomposition, aerosol material synthesis processes are also used to produce many other types of finely divided powders such as pigment grade titania, fumed silica, and carbon black [\[6](#page-6-0)]. In many of these cases, nucleation, surface reaction, coagulation and sintering have been studied in isolation, but how they interact in the free space of an aerosol tube reactor is not generally understood. Detailed knowledge of the fundamentals would allow industrial practitioners to develop new powder products and to produce powders with optimized properties.

To understand the details of this process, what is happening inside the black-box, we have built a model reactor that can be interrogated to provide detailed ''snap-shots'' of the particle population at many points during it's trajectory. For the first time ever, the details of the nickel and iron carbonyl powder making process can be studied and compared with models based on first principles. Our system consists of a 5 cm diameter tube that is 25 cm in height with five sampling arms perpendicular to the flow located axially every 5 cm. A photograph of the reactor is shown as Fig. 1. In this reactor, the carbonyl-laden feed gas is fed through a water-cooled nozzle into the top of the reactor. The furnace consists of five independent heating zones that conduct heat through the reactor wall. The

Fig. 1 Close-up of the 5 cm diameter tube reactor with the electrical heaters pulled away. On the left-hand side are the water jacketed arms and the pneumatic pistons that sample into the reactor

particles are collected in a filter bag at the bottom of the reactor.

For the experiment described here, three standard liters per minute of feedgas, comprised of 5 vol% nickel carbonyl and 330 ppm $NH₃$ in a balance of carbon monoxide was fed into the top of the reactor. The outside wall temperatures were maintained at about $600 °C$ and the system pressure was 1 psig. The internal temperatures, measured by in-situ temperature measurements are shown in Table 1. The top sample arm was deployed for 150 ms and the remaining four arms were each deployed for 1,000 ms. Two minutes elapsed between each deployment to allow the system to return to steady state.

Each sampling arm contains a sampler that has two flattened edges, which have been each fitted with three holders containing standard transmission electron microscope (TEM) grids. The foreground of Fig. [2](#page-2-0) shows one of the five sampling arms using in this study. The inset panel shows one of the carbon-coated copper TEM grids used on the sampler arms. The mechanism of particle deposition onto the grids is thermophoresis, this method has been used by a number of researchers elsewhere [\[7](#page-6-0), [8](#page-6-0)]. A thermophoretic force acts on the hot aerosol particles and pushes them towards cold surfaces [\[9](#page-6-0)]. This is an ideal method of

Table 1 The residence time and temperature, from CFD analysis, and in-situ sampling respectively, for the reactor centre-line position of each sampler

Sampler	Internal states within the reactor corresponding to sampler positions		
	Distance from inlet (cm)	Residence time from inlet (s)	Temperature $^{\circ}\mathrm{C}$
	5	0.01	350
2	10	0.08	410
3	15	0.37	430
	20	0.73	460
	25	1.11	450

Fig. 2 One of the thermophoretic samplers used in this study. Inset: a low magnification micrograph of one of the sample grids used in this study

capturing particles since it is essentially independent of particle size when the particles are smaller than the meanfree path of the gas, and when they are not, the decreased mobility can be estimated using well defined expressions such as Talbot's equation [[9\]](#page-6-0). The deployment of the samplers will now be described.

Each sampler resides in a sampler side-arm, shown in Fig. [1](#page-1-0). In the un-deployed state, the sample arm is isolated from the free space of the reactor. A small flow of inert gas is bled through the sample arm to insure that the reactor contents cannot come in contact with the sampler, except during deployment. Each sampler side-arm is also cooled by an annular water-cooled jacket, shown in Fig. [1](#page-1-0), which insures that there is an adequate temperature difference between the grids and the aerosol to drive the thermophoretic deposition. The objective was to operate the samplers in a regime where thermophoresis was at least ten times more significant than Brownian diffusion, thus avoiding a deposition mechanism that would biase the population [[9\]](#page-6-0). Each sample arm was coupled to a high speed double-action pneumatic piston. The timer/actuator for the pistons allowed for a precise deployment of each sampler, with approximately 20 ms of travel time into and out of the reactor. By this method, the sample grids could be inserted into the reactor parallel to the flow and ''exposed'' for precisely controlled periods of time.

The particle populations at the centre-line for each sampler are shown as Fig. [3](#page-3-0). Figure 3a show a scanning transmission electron microscope (STEM) (JEOL 2010F) micrograph of samples taken from the centre-line of the first sampler (5 cm from the inlet). This picture reveals the presence of a relatively monodisperse population of very fine particles, with an average size of 50 nm. No large particles were detected on the sample grid. Figure [3b](#page-3-0), c shows micrographs from the second sampler, observed at magnifications of $5,000 \times$ (SEM) and $60,000 \times$ (STEM), respectively. These images suggest that there is a bimodal population at the position of the second sampler. Figure [3d](#page-3-0), e shows low magnification images from the third and fourth sampler. High magnification images from these samplers, not shown here, reveal that there are no longer any particles less than 50 nm after the second sampler.

Computational fluid dynamics is a mathematical method for calculating the velocity, concentration and temperature profiles by dividing the geometry of interest into a number of adjacent cells and then solving discretized versions of the conservation equations for each cell. The technique has been used extensively to understand the temperature and flow patterns in aerosol reactors, assuming in most cases that the particle phase is sufficiently dilute to have any effect on the fluid states [\[10](#page-6-0)]. A detailed computational fluid dynamic model of this system was also developed to estimate the residence time of the gas at the centre-line; shown in Table [1](#page-1-0). Because of the small size of the particles in Fig. [3](#page-3-0), they will effectively track the bulk flow of the gas, in other words, they will have the same residence time distribution as the gas in this reactor.

Ten micrographs of the type shown in Figure [3](#page-3-0) were analysed by image analysis to derive population statistics such as number-weighted histograms for each of the 5 samplers at the centre-line. For segments of the histograms representing particles greater than the mean-free path of the gas, a correction was applied, using Talbot's equation to account for the decreased mobility (and hence observed frequency) of larger particles. Then these histograms were transformed to volume-weighted histograms by scaling each segment in the distribution by the corresponding average volume. Using standard statistical equations [\[11](#page-6-0)], the volume-weighted histograms could be used to calculate the geometric mean diameter (GMD) and geometric standard deviation (GSD) of these distributions, assuming that the population can be reasonably described by a log-normal relationship. In a log-normal distribution, the GMD is the average size and the GSD is a measure of the spread, or polydispersity of the distribution. The resulting picture of the evolving aerosol population for this experiment is shown as Fig. [4](#page-4-0).

In addition to estimates of the size distribution, estimates of the particle number concentration at each sampling position in the reactor can also be calculated. The thermophoretic velocity can be calculated using Talbot's equation, and the frequency of particles that have collected on the sample grids can also be measured. This data can then be combined to give estimates of the particle number concentration in the aerosol stream.

Figure [4](#page-4-0) shows some very interesting features that were corroborated by 20 other experiments using the same apparatus [[12\]](#page-6-0). The homogeneous reaction, or nucleation step occurs very early in the process, at a much lower

Fig. 3 Representative images from along the centre-line of the experiment. (a) the particles formed 5 cm from the inlet (residence time 0.01 s) by STEM. (b) The particles at 10 cm from the inlet (residence time 0.08 s) by SEM. (c) The particles at 10 cm from the

temperature than predicted by classical nucleation theory. Within 0.01 s, approximately 10^{15} particles per gram of gas appear that have consumed about 70% of the starting nickel carbonyl [\[12\]](#page-6-0). High resolution TEM confirms that these particles are single crystals, suggesting that these particles have grown mainly by a surface reaction, and not coagulation, which would have produced randomly oriented poly-crystals.

The most interesting behaviour of the system occurs between 0.01 and 0.08 s. During this time, the 50 nm particles have begun to combine by coagulation, this is expected to happen very quickly since Brownian coagulation is proportional to the square of the number concentration. The nucleation reaction has not been completely disabled, so that smaller particles concurrently enter the distribution from the left-hand side of the histogram in Fig. [4](#page-4-0). They grow more slowly than the particles observed at the first sampler because of the low nickel carbonyl concentration, resulting in a finer size. The simultaneous action of nucleation, surface reaction and coagulation have produced a transient bimodal distribution with a maximum polydispersity of GSD = 3.90. The particle concentration is

inlet under high magnification by STEM. (d) The particles at 15 cm from the inlet (residence time 0.37 s) by SEM. (e) The particle at 20 cm from the inlet (residence time 0.73 s) by SEM

still extremely high, approximately 10^{15} – 10^{16} particles per gram of gas.

By the position of the third sample, corresponding to a residence time of 0.37 s, nucleation has been extinguished and the particles continue to grow mainly by Brownian coagulation. The particle concentration has decayed to approximately 10^{11} – 10^{12} particle per gram of gas. As the particles continue down through the reactor without any new reactant, they continue to grow by collisions, and the spread of the distribution approaches a constant value, which is expected for a coagulation-dominated process [\[9](#page-6-0)].

This fascinating picture of the aerosol growth process has been uncovered experimentally for the first time by these experiments. This phenomenon was previously predicted mathematically for the oxidation of $TiCl₄$ in an aerosol tube reactor using a sectional aerosol model [[13,](#page-6-0) [14](#page-6-0)]. Both of these groups showed that the polydispersity starts off at a minimum as the reaction begins by nucleation and surface reaction, then goes through a maximum as a result of a temporary bimodal distribution, when nucleation, surface reaction, and coagulation all co-exist. During the temporary bimodal distribution, particles being Fig. 4 Mass weighted histograms, based on image analysis of micrographs collected from the samplers. After 0.01 s, the distribution is narrow with all particles less than about 100 nm. At 0.08 s, the distribution has become strongly bimodal as nucleation, surface reaction and coagulation all co-exist. Some time after 0.08 s, nucleation is extinguished and particle growth continues by the surface reaction and coagulation. As coagulation dominates, the spread of the distribution approaches the self-preserving limit

0 0.02 0.04 0.06 0.08 0.1 0.3 0.5 0.7 0.9 1.1 1.3 1.5 1.7 1.9 2.1 2.3 2.5 2.7 2.9

Particle size (µ**m)**

produced at the fine end of the distribution are smaller than particles produced during the initial stages of the process. Finally the population approaches the limiting value of the polydispersity as the distribution approaches the self-preserving limit, which is characteristic for aerosol systems growing by Brownian coagulation only. To the best knowledge of the authors, the present work is the first experimental corroboration of this numerically predicted phenomenon.

0.1

0 0.05 0.1 0.15

 Ω 0.05 0.1 0.15

Mass fraction in size segment (fraction/µ**m)**

Mass fraction in size segment (fraction/um) 0.2

 $\overline{0}$ 0.1

The consequences of these observations can be significant for practitioners trying to make metal powders with optimized properties. For example, in many applications, powders with a narrow size distribution are preferred. According to our results, the process is dominated by Brownian coagulation after 0.1 s. After this point, the minimum polydispersity or GSD of the population is fundamentally limited as the ''self-preserving size distribution'' is approached. Of course, the actual polydispersity can be higher, depending on the degree of fluid mixing in the reactor.

One of the interesting implications of this view within the reactor, is that sintering of the particles must be very rapid, in fact the particles that come out of the bottom of the reactor, Fig. [3](#page-3-0)e do not seem to be comprised of the fine particles shown in Fig. [3a](#page-3-0). To prove that this can be accounted for by coagulation and sintering, the time constants for particle sintering by a variety of mechanisms were calculated for nickel using the 2-sphere model [\[15](#page-6-0)]. This model provides an estimate for the sintering time of two equi-sized particles. Many of the initial particles

 $GMD = 0.99 \,\mu m$ **GSD = 1.96**

ninnn

consist of platelets and should sinter very effectively. At the prevailing internal temperature profile of this study, approximately 440 \degree C and 1 s, the maximum particle size for initial sintering by surface diffusion, evaporation–condensation, grain boundary diffusion and volume diffusion is 0.32, 0.005, 0.17 and $< 0.001 \,\mu m$ [\[12](#page-6-0)]. Because the largest quasi-spherical particles are approximately $0.7 \mu m$, the conclusion is that the predominant mechanism for sintering under these conditions is surface diffusion, the temperature is too high for significant evaporation–condensation via a carbonyl intermediary and it is too low for grain-boundary or volume diffusion.

High resolution field emission scanning electron microscope (SEM) images of typical particles at each sampler suggest that the particles have grown by coagulation and sintering, shown in Fig. 5. Figure 5a shows the plate-like morphology of many of the very fine particles produced after 0.01 s. Figure 5b, a typical particle after 0.08 s suggests that a 0.5 micron particle has grown by the accumulation of many of the smaller particles present earlier. Further displacement through the hot reactor shows a progressive smoothing of the particle surface (5c corresponds to 0.37 s and 5d corresponds to 0.73 s). Near the bottom of the reactor, the surface has been sufficiently smoothed that it is no longer obvious how the particle has grown.

As final proof of this mechanism, a focused-ion beam was used to section a typical particle. The particle was subsequently imaged in back scattered electron (BSE) mode to contrast the crystallographic orientation inside the particle—this image is shown in Fig. 6. The exterior of the

Fig. 6 A focused ion beam section of a nickel particle made from nickel carbonyl, revealing details of the internal crystal structure

Fig. 5 Typical particles at the centre-line as viewed by high resolution field emission SEM. (a) At 5 cm (0.01 s), nucleation is much faster than coagulation. (b) At 10 cm (0.08 s) smaller particles

have collided together and melded by surface diffusion. (c) (15 cm, 0.37 s), (d) (20 cm, 0.73 s) as the residence time increases, samples show less detail that reveal their genesis

particle shown in Fig. [6](#page-5-0) is smooth, suggesting significant surface sintering, however on the inside, there are crystals, approximately 100 nm, the same order of magnitude as the fine particles produced in the top of the reactor (Fig. [3a](#page-3-0)). There is no observed orientation relationship between adjacent crystals inside the particle, which is consistent with growth by coagulation.

Specific engineering of the material properties of aerosol-produced powders becomes a possibility when each step of the process is understood from the perspective of the scientific fundamentals. This deep process knowledge has become available for the metal carbonyl systems through the combination of in-situ sampling and CFD analysis, the ''black-box'' has finally been opened up. For the temperature ranges used in this process, the crystal size is determined by nucleation and surface reaction, which occurs very early in the overall process. The smoothness of the particle surface is controlled by surface sintering, evaporation–condensation below about 300 °C or surface diffusion above this temperature. Finally, the size and shape of the size distribution are controlled by coagulation. The spread of the distribution has a lower limit, which can be calculated from aerosol physics. Finally, the system was shown to behave similarly to mathematical models developed for other aerosol synthesis processes. It is expected that many of these insights will carry over to other aerosol synthesis systems.

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References

- 1. Mond R (1930) J Soc Chem Ind XLIX(24):271, (25):283, (26):287
- 2. Barin I (1989) Thermochemical data of pure substances VCH
- 3. Trout W (1938) J Chem Educ March:113
- 4. www.incosp.com
- 5. Carlton H, Oxley J (1965) Am Inst Chem Eng J 11(1): 79
- 6. Pratsinis S, Vemury S (1996) Powd Technol 88:267
- 7. Dobbins R, Megaridis C (1987) Langmuir 3:254
- 8. Arabi-Katbi O, Pratsinis S, Morrison P, Megaridis C (2001) Combust Flame 124:560
- 9. Friedlander S (2000) Smoke, dust and haze, 2nd edn. John Wiley
- 10. Kammler H, Madler L, Pratsinis S (2001) Chem Eng Technol 24:583
- 11. Hinds W (1999) Aerosol technology: properties, behavior, and measurement of airborne particles, 2nd edn. John Wiley and Sons
- 12. Wasmund E (2005) Powder making in an aerosol tube reactor. PhD, Department of Materials Science and Engineering, McMaster University
- 13. Xiong Y, Pratsinis S (1991) J Aerosol Sci 22(5):637
- 14. Tsantilis S, Pratsinis S (2000) Am Inst Chem Eng J 46(2):407
- 15. German R (1996) Sintering theory and practice. John Wiley